

and the diene **3** to predict the difference, ΔE_A , between the E_A 's of ethylene and the monoene **2**. This procedure yields $x = 0.27$ eV and $y = 0.14$ eV, or $x - y = +0.13$ eV, in agreement with experiment. The σ^*,π^* and σ',π^* interactions are significantly smaller than the σ,π interactions in the same systems on account of the large energy gap between the $\pi^* - \pi^*$ MO and the σ' MO and the relatively poor overlap between the $\pi^* + \pi^*$ MO and the σ^* MO.

The present study indicates two salient features of π^* -OITB: (1) these interactions result from the mixing of the $\pi^* \pm \pi^*$ symmetry-adapted orbitals with both σ and σ^* frontier MO's, with the latter predominating (i.e., $x > y$); (2) the net splitting in the π^* manifolds is nearly the same as in the π manifolds. This follows because the $\pi^* \pm \pi^*$ combinations are shifted in *opposite* directions through mixing with the σ and σ^* MO's, while the $\pi \pm \pi$ combinations are shifted in the *same* direction through mixing with the occupied σ orbitals.

Acknowledgment. This research was supported by the National Science Foundation (P.D.B. and K.D.J.) and the Australian Research Grants Scheme (M.N.P.R.). We thank K. N. Houk for valuable discussions and L. Ng for assistance with the experiments.

Registry No. **2**, 1076-12-6; **3**, 83602-18-0; **4**, 36012-52-9; **5**, 15914-93-9; **6**, 15914-94-0; **7**, 1076-12-6; **8**, 1076-13-7; norbornene, 498-66-8.

Diprotonated Methane, CH_6^{2+} , and Diprotonated Ethane, $\text{C}_2\text{H}_8^{2+}$

Koop Lammertsma* and George A. Olah

Hydrocarbon Research Institute
and Department of Chemistry
University of Southern California
Los Angeles, California 90089-1662

Mario Barzagli and Massimo Simonetta

Institute of Physical Chemistry and CNR Center
University of Milan, 20133 Milan, Italy

Received July 2, 1982

The stabilization of the *pentacoordinate* CH_5^+ ion¹ finds its origin in the interaction of the electron-deficient carbon with two ligands (three center, two electron (3c-2e)). The size of this effect is reflected in the heat of hydrogenation of CH_3^+ , experimentally deduced as 40 kcal/mol^{2a} and calculated as 34.5 kcal/mol by ab initio molecular orbital theory (MP3/6-31G** + ZPE).^{2b} The 3c-2e stabilization is a common phenomenon in larger systems, like the norbornyl cation,³ although its magnitude reduces as illustrated by the heats of hydrogenation of C_2H_5^+ (11.8 kcal/mol) and C_3H_7^+ (2.3 kcal/mol).^{2a}

The concept of hypercoordination of electron-deficient species can also be applied to dications. From our ab initio theoretical⁴ studies on small carbocations, the prototypes of new unprecedented types of hypercoordinate ions emerged, i.e., diprotonated methane,

Table I. Total (hartrees) and Relative Energies (kcal/mol) Using the 6-31G** Basis and HF/6-31G* Geometries and Zero-Point Energies (kcal/mol)

geometry	HF	MP3	zero point
1, C_{2v}	-40.30397 (0)	-40.49674 (0)	37.9
2, C_{2v}	-40.29730 (4.2)	-40.49264 (2.6)	35.8
3, C_2	-79.45620 (0)	-79.80097 (0)	58.1 ^a
4, D_{4d}	-79.44122 (9.4)	-79.79618 (3.0)	55.8 ^a
7, C_{2v}	-78.27287 (5.9)	-78.56410 (3.5)	43.2 ^a
8, C_s	-78.28234 (0)	-78.56969 (0)	44.4 ^a
5, D_{2d}	-77.09538	-77.32626	30.1 ^a

^a Zero-point vibrational energy obtained with the 3-21G basis.

Scheme I

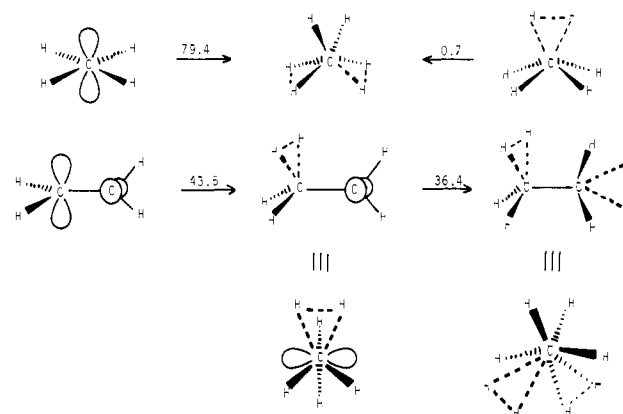
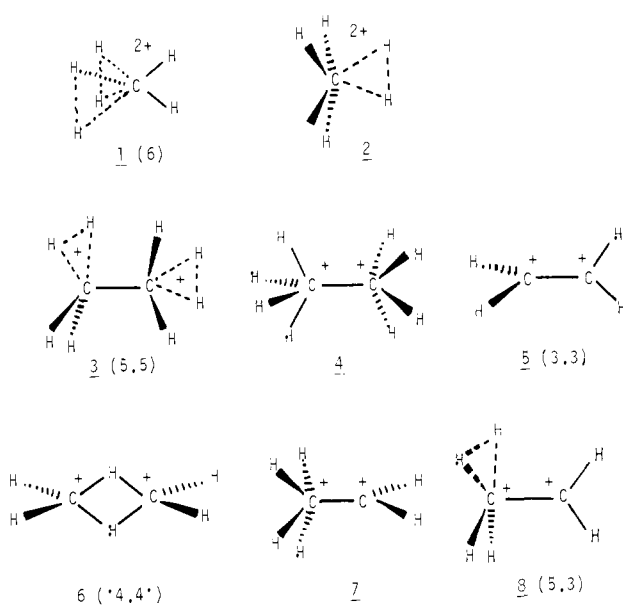


Chart I



CH_6^{2+} , and diprotonated ethane, $\text{C}_2\text{H}_8^{2+}$. In the context of pentacoordinate carbon stabilization we also restudied the ethane dication, $\text{C}_2\text{H}_6^{2+}$.

In order to describe these hypercoordinate species, we have used single-point calculations with the flexible 6-31G** basis set, which has polarization functions on both hydrogen and carbon, for geometries optimized within the Hartree-Fock limit at the 6-31G* (d functions) level.⁵ Furthermore, at this level we have included electron correlation corrections with the Møller-Plesset perturbation theory^{6,7} to the third order with frozen core (MP3/6-

(1) Olah, G. A.; Klopman, G. A.; Schlosberg, R. H. *J. Am. Chem. Soc.* **1969**, *91*, 3261.

(2) (a) Hiraoka, K.; Kebarle, P. *J. Am. Chem. Soc.* **1976**, *98*, 6119. (b) Raghavachari, K.; Whiteside, R. A.; Pople, J. A.; Schleyer, P. v. R. *Ibid.* **1981**, *103*, 5649.

(3) (a) Olah, G. A.; Prakash, G. K. S.; Arvanaghi, M.; Anet, F. A. L. *J. Am. Chem. Soc.*, in press. (b) Olah, G. A. *Acc. Chem. Res.* **1976**, *9*, 41. (c) Goddard, J. D.; Osamura, Y.; Schaefer H., III *J. Am. Chem. Soc.* **1982**, *104*, 3258.

(4) A slightly modified version of the Gaussian 80 series of programs (see ref 4a) for IBM computers was used. (a) Binkley, J. S.; Whiteside, R. A.; Krishnan, R.; Seeger, R.; DeFrees, D. J.; Schlegel, H. B.; Topiol, S.; Kahn, L. R.; Pople, J. A. *QCPE* **1981**, *13*, 406. (b) Van Kampen, P. N.; Smits, G. F.; De Leeuw, F. A. A. M.; Altona, C. *Ibid.* **1982**, *14*, 437.

(5) For the 6-31G* and 6-31G** basis sets, see: Hariharan, P. C.; Pople, J. A. *Theor. Chim. Acta* **1973**, *28*, 213. For the 3-21G basis set, see: Binkley, J. S.; Pople, J. A.; Hehre, W. J. *J. Am. Chem. Soc.* **1980**, *102*, 939. Gordon, M. S.; Binkley, J. S.; Pople, J. A.; Pietro, W. J.; Hehre, W. J. *Ibid.* **1982**, *104*, 2797.

31G**) for all data and used these throughout the text (see Table I). All species reported are considered to be equilibrium structures, since the force-constant matrix obtained from the 6-31G* (CH₆²⁺) and 3-21G⁵ (C₂H₄²⁺, C₂H₆²⁺, and C₂H₈²⁺) optimized geometries have no negative eigenvalues, except **2**, which has one. The zero-point vibrational energies deduced from these force constants (see Table I) are after scaling⁸ included as our final adjustment in all subsequent reaction energies given at the MP3/6-31G** level (Scheme I).

The methane dication, CH₄²⁺, is trivalent and tetracoordinate with a planar (*D*_{4h} symmetry) geometry.⁹ Interaction of molecular hydrogen with the empty p_z orbital of CH₄²⁺ results in CH₆²⁺. Diprotonated methane is *hexacoordinate*. The minimum-energy equilibrium structure **1** (Chart I) calculated for CH₆²⁺ (see Table I) has two orthogonal 3c-2e interactions (*C*_{2v} symmetry), emphasizing the importance of this effect. The structure **2** with only one such interaction may be regarded as the transition for transfer of one 3c-2e interaction in **1** and is 2.6 kcal/mol higher in energy. The stabilization of **1** gained with respect to CH₄²⁺ amounts to 79.4 kcal/mol. Intuitively one expects carbocations to be highly unstable because of electrostatic repulsion. Indeed CH₆²⁺ has a high exothermicity of 63.1 kcal/mol for proton loss and of 126.8 kcal/mol toward dissociation to CH₃⁺ and H₃⁺. However, our studies indicate a barrier of 40 kcal/mol (MP3/6-31G**) for the deprotonation and suggest even more for loss of H₃⁺. Since CH²⁺, CH₂²⁺, CH₃²⁺, and CH₄²⁺ have already been reported in a gas-phase study,¹⁰ CH₆²⁺ may be also a viable species. The calculated heat of formation of 651 ± 3 kcal/mol is actually the lowest of the C₁ dications.

In diprotonated ethane, C₂H₆²⁺, each carbon is *pentacoordinate*. We may regard C₂H₈²⁺ as hydrogenated ethane dication, C₂H₆²⁺ (see later). Even better, C₂H₈²⁺ can be considered as doubly hydrogenated ethylene dication, C₂H₄²⁺, with each of the orthogonal vacant p_z orbitals in strong interaction with a hydrogen molecule. This view is confirmed through the calculations. For the equilibrium structure with minimum energy we find C₂ symmetry (**3**) with the two 3c-2e interactions at about a 90° dihedral angle. The Newman projection (Scheme I) shows that complexation of the two hydrogen molecules only slightly distorts the perpendicular ethylene dication skeleton **5** (*D*_{2d} symmetry). At the 6-31G** level **3** is 9.4 (HF) and 3.0 (MP3) kcal/mol more stable than the second equilibrium structure **4** (*D*_{4d} symmetry) and therefore underlines the importance of 3c-2e interactions. The heat of hydrogenation of C₂H₄²⁺ to give diprotonated ethane is 79.8 kcal/mol. The most likely dissociation of C₂H₈²⁺ is toward CH₅⁺ and CH₃⁺ and is exothermic by 100.1 kcal/mol. A significant barrier for dissociation is expected, however, because of the distinct reorganization of atoms needed in the transition state.

We reported recently¹¹ the structure of the ethane dication as the doubly bridged, diborane-like dication **6** (*D*_{2d} symmetry). Schleyer et al.¹² subsequently found the carbenium-carbonium structure **7** (*C*_{2v} symmetry) to be 9.0 kcal/mol more stable (MP4SDQ/6-31G**//6-31G*). The preference of **7** was explained by the reduced electrostatic repulsion for the hydrogens lying farther apart and the stronger C-C bonding through hyperconjugation.¹² Structure **7** has a *tri- and a tetravalent* carbon with *tri- and pentacoordination*, respectively. In the context of the present concept one might formulate C₂H₆²⁺ as hydrogenated

C₂H₄²⁺. Complexation of molecular hydrogen with a vacant orbital on one of the carbons of the ethylene dication **5** is then expected to result in the carbenium-carbonium dication C₂H₆²⁺, with the carbonium center involved in a 3c-2e interaction as in **8** (*C*_s symmetry). This is indeed confirmed by our calculations on a reinvestigation of C₂H₆²⁺. The equilibrium structure **8** is 3.5 kcal/mol (MP3/6-31G**) more stable than **7**, reported by Schleyer et al.¹² (see Table I). The Newman projection of **8** shows the orthogonality of the 3c-2e interaction and the vacant p orbital and also suggests a somewhat larger hyperconjugation as in **7**. The hydrogenation of the ethylene dication **5** is exothermic by 43.5 kcal/mol. The kinetic stability of C₂H₄²⁺ combined with the reported barrier for dissociation of C₂H₆²⁺,¹² now adjusted for **8**, of 30 kcal/mol may render the ethane dication feasible for experimental observation.

The calculations further suggest that all three hypercoordinate dications, CH₆²⁺, C₂H₈²⁺, and C₂H₆²⁺, with coordination numbers 6, 5-5, and 5-3, respectively, may be viable species. A possible route to experimental observation could be via hydrogenation of the appropriate carbocationic precursors, obtained by the mass spectroscopic charge-stripping technique in the presence of hydrogen in the neutral gas.

Acknowledgment. The cooperation of the USC computer center greatly facilitated the present investigation. M.B. thanks the Italian Consiglio Nazionale delle Ricerche for the award of a National Scholarship for research abroad.

Registry No. CH₆²⁺, 83561-00-6; C₂H₈²⁺, 83561-01-7; CH₄²⁺, 34557-54-5; C₂H₄²⁺, 54509-73-8.

Bimolecular Substitution at Carbon in Neopentyl-Like Silylcarbonyl Sulfonates

Peter J. Stang* and Mladen Ladika

Department of Chemistry, University of Utah
Salt Lake City, Utah 84112

Yitzhak Apeloig* and Amnon Stanger

Department of Chemistry
Technion—Israel Institute of Technology
Technion City, Haifa 32000, Israel

Melvyn D. Schiavelli* and Michael R. Hughey

Department of Chemistry, College of William and Mary
Williamsburg, Virginia 23185

Received August 16, 1982

There is considerable current synthetic and mechanistic interest in silicon chemistry. It has been reasonably well established that carbocations are stabilized by β-silicon substitution and largely by implication that they are destabilized by α-silicon substitution.^{1,2} For example, no detectable reaction of (CH₃)₃SiCH₂X (X = Br, Cl) was observed in aqueous solvents at 70 °C.³ Furthermore, PhCMe(SiMe₃)Br reacts slower than PhC(CH₃)₂Br, as does (CH₃)₃SiC(CH₃)₂Br compared to (CH₃)₃CC(CH₃)₂Br.⁴

The effect of silicon substitution at the reaction center on a bimolecular displacement at carbon is even less well understood. Thus, while (CH₃)₃SiCH₂Br reacts 2600 times more rapidly with ethoxide in ethanol than does (CH₃)₃CCH₂Br,⁵ (CH₃)₃SiCH₂Cl reacts slower than CH₃CH₂Cl with I⁻ in aqueous ethanol and faster than CH₃CH₂Cl with I⁻ in acetone.³

The solvolysis reaction of neopentyl-X compounds proceeds with considerable *k*_A character in many solvents, resulting in

- (1) Lambert, J. B.; Finzel, R. B. *J. Am. Chem. Soc.* **1982**, *104*, 2020-2022.
- (2) Brook, A. G.; Pannell, K. H. *Can. J. Chem.* **1970**, *48*, 3679-3693.
- (3) Cook, M. A.; Eaborn, C.; Walton, D. R. M. *J. Organomet. Chem.* **1971**, *29*, 389-396.
- (4) Cartledge, F. K.; Jones, J. P. *Tetrahedron Lett.* **1971**, 2193-2196.
- (5) Dostrovsky, I.; Hughes, E. D. *J. Chem. Soc.* **1946**, 157-161.

(6) (a) Möller, C.; Plesset, M. S. *Phys. Rev.* **1934**, *46*, 618. (b) Pople, J. A.; Binkley, J. S.; Seeger, R. *Int. J. Quantum Chem. Symp.* **1976**, *10*, 1.

(7) For the influence of electron correlation by Möller-Plesset perturbation theory see ref 2b. Also see: Bartlett, R. J. *Ann. Rev. Phys. Chem.* **1981**, *32*, 359.

(8) The zero-point energies are scaled by 0.9; see: Pople, J. A.; Schlegel, H. B.; Krishnan, R.; DeFrees, D. J.; Binkley, J. S.; Frisch, M. J.; Whiteside, R. A.; Hout, R. F.; Hehre, W. J. *Int. J. Quantum Chem. Symp.* **1982**, *15*, 269.

(9) Pople, J. A.; Tidor, B.; Schleyer, P. v. R. *Chem. Phys. Lett.* **1982**, *88*, 533. Siegbahn, P. E. M. *Chem. Phys.* **1982**, *66*, 443.

(10) Ast, T.; Porter, C. J.; Proctor, C. J.; Beynon, J. H. *Chem. Phys. Lett.* **1981**, *78*, 439.

(11) Olah, G. A.; Simonetta, M. *J. Am. Chem. Soc.* **1982**, *104*, 330.

(12) Schleyer, P. v. R.; Kos, A. J.; Pople, J. A.; Balaban, A. T. *J. Am. Chem. Soc.* **1982**, *104*, 3771.